

Time Dependent Density Functional Theory

An Introduction

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Belfast, 29 Jun 2007



Outline

- 1 Introduction: why TD-DFT ?
- 2 (Just) A bit of Formalism
 - TDDFT: the Foundation
 - Linear Response Formalism
- 3 TDDFT in practice:
 - The ALDA: Achievements and Shortcomings
- 4 Resources

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Density Functional ... Why ?

Basic ideas of DFT

- 1 Any observable of a quantum system can be obtained from the density of the system **alone**.
- 2 The density of an interacting-particles system can be calculated as the density of an auxiliary system of **non-interacting** particles.

Importance of the density

Example: atom of Nitrogen (7 electron)

$\Psi(\mathbf{r}_1, \dots, \mathbf{r}_7)$ 21 coordinates

10 entries/coordinate $\Rightarrow 10^{21}$ entries

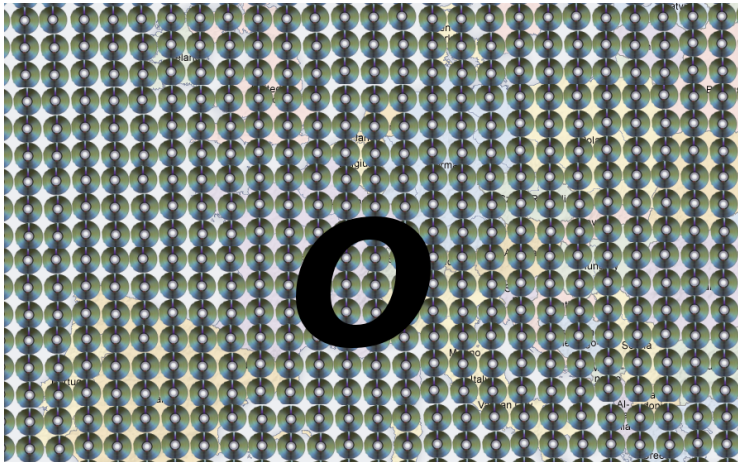
8 bytes/entry $\Rightarrow 8 \cdot 10^{21}$ bytes

4.7×10^9 bytes/DVD $\Rightarrow 2 \times 10^{12}$ DVDs

Density Functional ... Why ?



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Importance of the density

Example: atom of Carbon (6 electron)

$\Psi(\mathbf{r}_1, \dots, \mathbf{r}_6)$ 18 coordinates

10 entries/coordinate $\Rightarrow 10^{18}$ entries

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Importance of non-interacting

The Kohn-Sham one-particle equations

$$H_i(\mathbf{r})\psi_i(\mathbf{r}) = \epsilon_i(\mathbf{r})\psi_i(\mathbf{r})$$

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Density Functional ... Successfull ?

TABLE II: Top-10 cited PR articles. The asterisks denote citation undercount due to citations with missing prepended A/B page numbers – 123 out of 3227 total for item 1 and 120 out of 2640 for item 2.

Impact Rank	Publication			# cites	Av. Age	Impact	Title	Author(s)	
1	PR	140	A1133	1965	3227*	26.64	85972	Self-Consistent Equations...	W. Kohn & L. J. Sham
2	PR	136	B864	1964	2460*	28.70	70604	Inhomogeneous Electron Gas	P. Hohenberg & W. Kohn
3	PRB	23	5048	1981	2079	14.38	29896	Self-Interaction Correction to...	J. P. Perdew & A. Zunger
4	PRL	45	566	1980	1781	15.42	27463	Ground State of the Electron ...	D. M. Ceperley & B. J. Alder
5	PR	108	1175	1957	1364	20.18	27526	Theory of Superconductivity	J. Bardeen, L. N. Cooper, & J. R. Schrieffer
6	PRL	19	1264	1967	1306	15.46	20191	A Model of Leptons	S. Weinberg
7	PRB	12	3060	1975	1259	18.35	23103	Linear Methods in Band Theory	O. K. Andersen
8	PR	124	1866	1961	1178	27.97	32949	Effects of Configuration...	U. Fano
8	RMP	57	287	1985	1055	9.17	9674	Disordered Electronic Systems	P. A. Lee & T. V. Ramakrishnan
9	RMP	54	437	1982	1045	10.82	11307	Electronic Properties of...	T. Ando, A. B. Fowler, & F. Stern
10	PRB	13	5188	1976	1023	20.75	21227	Special Points for Brillouin-...	H. J. Monkhorst & J. D. Pack



S. Redner <http://arxiv.org/abs/physics/0407137>

Time Dependent DFT ... Why ?

Large field of research concerned with
many-electron systems in time-dependent fields

Different Phenomena

- absorption spectra
- energy loss spectra
- photo-ionization
- high-harmonic generation
- photo-emission

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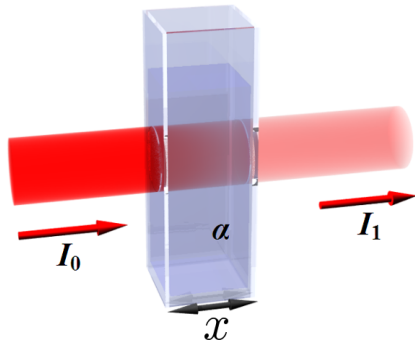
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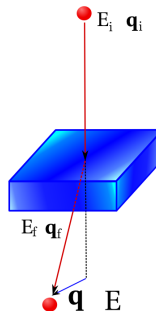


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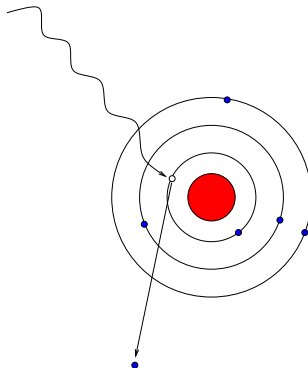


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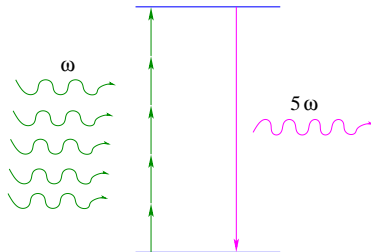


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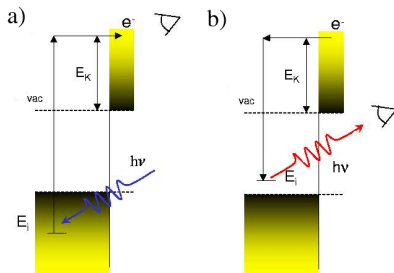


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Time Dependent DFT ... Why ?

We need a time dependent theory



TDDFT is a promising candidate

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The name of the game: TDDFT

DFT

Hohenberg-Kohn theorem 1

The ground-state expectation value of any physical observable of a many-electrons system is a unique functional of the electron density $n(\mathbf{r})$

$$\langle \varphi^0 | \hat{O} | \varphi^0 \rangle = O[n]$$



P. Hohenberg and W. Kohn

Phys.Rev. **136**, B864 (1964)

(Fermi, Slater)

TDDFT

Runge-Gross theorem

The expectation value of any physical time-dependent observable of a many-electrons system is a unique functional of the **time-dependent** electron density $n(\mathbf{r}, t)$ and of **the initial state**

$$\varphi^0 = \varphi(t=0)$$

$$\langle \varphi(t) | \hat{O}(t) | \varphi(t) \rangle = O[n, \varphi^0](t)$$



E. Runge and E.K.U. Gross

Phys.Rev.Lett. **52**, 997 (1984)

(Ando, Zangwill and Soven)

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DFT

Static problem

Second-order differential
equation

Boundary-value problem.

$$H\varphi(\mathbf{r}_1, \dots, \mathbf{r}_N) = E\varphi(\mathbf{r}_1, \dots, \mathbf{r}_N)$$

TDDFT

Time-dependent problem

First-order differential equation

Initial-value problem

$$H(t)\varphi(\mathbf{r}_1, \dots, \mathbf{r}_N; t) = i\hbar \frac{\partial}{\partial t} \varphi(\mathbf{r}_1, \dots, \mathbf{r}_N; t)$$

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- $V_{\text{ext}}(\mathbf{r}, t) \neq V'_{\text{ext}}(\mathbf{r}, t) \iff \mathbf{j}(\mathbf{r}, t) \neq \mathbf{j}'(\mathbf{r}, t)$
 - $\nabla \cdot [n \nabla V_{\text{ext}}] \neq \nabla \cdot [n \nabla V'_{\text{ext}}] \iff n(\mathbf{r}, t) \neq n'(\mathbf{r}, t)$
- $$n(\mathbf{r}, t) \longrightarrow V_{\text{ext}}(\mathbf{r}, t) + c(t) \longrightarrow \varphi e^{ic(t)}$$

$$\langle \varphi(t) | \hat{O}(t) | \varphi(t) \rangle = O[n, \varphi^0](t)$$

What about infinite systems?

The name of the game: TDDFT

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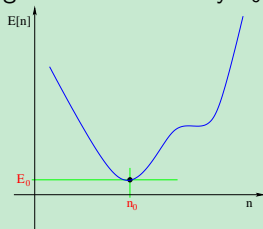
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DFT

Hohenberg-Kohn theorem 2

The total energy functional has a minimum, the ground-state energy E_0 , corresponding to the ground-state density n_0 .



TDDFT

Runge-Gross theorem - No minimum

Time-dependent Schrödinger eq. (initial condition $\varphi(t=0) = \varphi_0$), corresponds to a **stationary point** of the Hamiltonian action

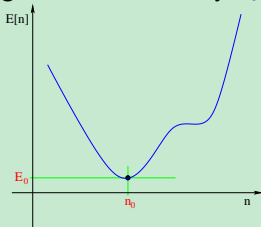
$$A = \int_{t_0}^{t_1} dt \langle \varphi(t) | i \frac{\partial}{\partial t} - H(t) | \varphi(t) \rangle$$

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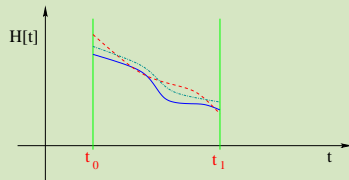


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Kohn-Sham equations

$$\left[-\frac{1}{2} \cdot \nabla_i^2 + V_{\text{tot}}(\mathbf{r}) \right] \phi_i(\mathbf{r}) = \epsilon_i \phi_i(\mathbf{r})$$

$$V_{\text{tot}}(\mathbf{r}) = V_{\text{ext}}(\mathbf{r}) + \int d\mathbf{r}' v(\mathbf{r}, \mathbf{r}') n(\mathbf{r}') + V_{\text{xc}}([n], \mathbf{r})$$

$$V_{\text{xc}}([n], \mathbf{r}) = \frac{\delta E_{\text{xc}}[n]}{\delta n(\mathbf{r})}$$

Unknown exchange-correlation potential.

V_{xc} functional of the density.

TDDFT

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The name of the game: TDDFT

Demonstrations, further readings, etc.



R. van Leeuwen

Int.J.Mod.Phys. **B15**, 1969 (2001)

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Causality-Symmetry dilemma

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First Approach: Time Evolution of KS equations

$$[H_{KS}(t)] \phi_i(\mathbf{r}, t) = i \frac{\partial}{\partial t} \phi_i(\mathbf{r}, t)$$

$$n(\mathbf{r}, t) = \sum_i^{\text{occ}} |\phi_i(\mathbf{r}, t)|^2$$

$$\phi(t) = \hat{U}(t, t_0) \phi(t_0)$$

$$U(t, t_0) = 1 - i \int_{t_0}^t d\tau H(\tau) \hat{U}(\tau, t_0)$$



A. Castro *et al.* J.Chem.Phys. **121**, 3425 (2004)

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First Approach: Time Evolution of KS equations

Photo-absorption cross section σ

$$\sigma(\omega) = \frac{4\pi\omega}{c} \text{Im}\alpha(\omega)$$

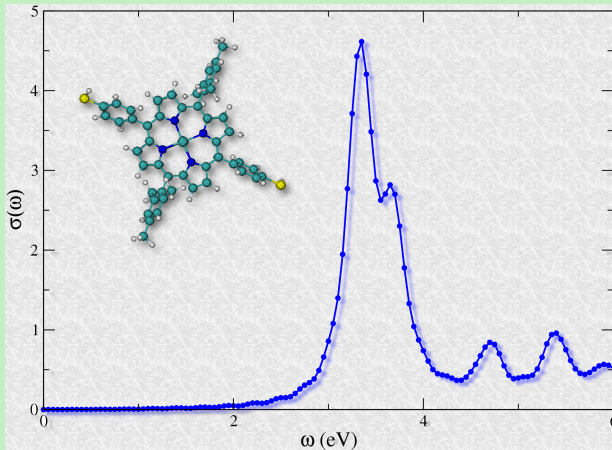
$$\alpha(t) = - \int d\mathbf{r} V_{\text{ext}}(\mathbf{r}, t) n(\mathbf{r}, t)$$

in dipole approximation ($\lambda \gg \gg$ dimension of the system)

$$\sigma_{zz}(\omega) = -\frac{4\pi\omega}{c} \text{Im} \alpha(\omega) = -\frac{4\pi\omega}{c} \text{Im} \int d\mathbf{r} z n(\mathbf{r}, \omega)$$

First Approach: Time Evolution of KS equations

Photo-absorption cross section σ : porphyrin



First Approach: Time Evolution of KS equations

Other observables

Multipoles

$$M_{lm}(t) = \int d\mathbf{r} r^l Y_{lm}(r) n(\mathbf{r}, t)$$

Angular momentum

$$L_z(t) = - \sum_i \int d\mathbf{r} \phi_i(\mathbf{r}, t) i (\mathbf{r} \times \nabla)_z \phi_i(\mathbf{r}, t)$$

First Approach: Time Evolution of KS equations

Advantages

- Direct application of KS equations
- Advantageous scaling
- Optimal scheme for finite systems
- All orders automatically included

Shortcomings

- Difficulties in approximating the $V_{xc}[n](\mathbf{r}, t)$ functional of the history of the density
- Real space not necessarily suitable for solids
- Does not explicitly take into account a “small” perturbation. Interesting quantities (excitation energies) are contained in the linear response function!

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Linear Response Approach

Polarizability

interacting system $\delta n = \chi \delta V_{ext}$

non-interacting system $\delta n_{n-i} = \chi^0 \delta V_{tot}$

Linear Response Approach

Polarizability

interacting system $\delta n = \chi \delta V_{ext}$

non-interacting system $\delta n_{n-i} = \chi^0 \delta V_{tot}$

Single-particle polarizability

$$\chi^0 = \sum_{ij} \frac{\phi_i(\mathbf{r})\phi_j^*(\mathbf{r})\phi_i^*(\mathbf{r}')\phi_j(\mathbf{r}')}{\omega - (\epsilon_i - \epsilon_j)}$$

hartree, hartree-fock, dft, etc.



G.D. Mahan *Many Particle Physics* (Plenum, New York, 1990)

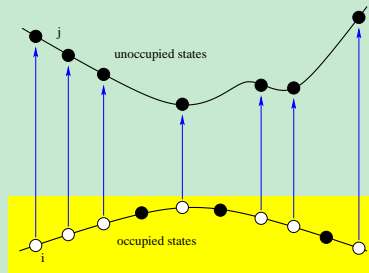
Linear Response Approach

Polarizability

interacting system $\delta n = \chi \delta V_{ext}$

non-interacting system $\delta n_{n-i} = \chi^0 \delta V_{tot}$

$$\chi^0 = \sum_{ij} \frac{\phi_i(\mathbf{r})\phi_j^*(\mathbf{r})\phi_i^*(\mathbf{r}')\phi_j(\mathbf{r}')}{\omega - (\epsilon_i - \epsilon_j)}$$



Linear Response Approach

Polarizability

interacting system $\delta n = \chi \delta V_{\text{ext}}$

non-interacting system $\delta n_{n-i} = \chi^0 \delta V_{\text{tot}}$



Density Functional Formalism

$$\delta n = \delta n_{n-i}$$

$$\delta V_{\text{tot}} = \delta V_{\text{ext}} + \delta V_H + \delta V_{\text{xc}}$$

Linear Response Approach

Polarizability

$$\chi \delta V_{\text{ext}} = \chi^0 (\delta V_{\text{ext}} + \delta V_H + \delta V_{\text{xc}})$$

$$\chi = \chi^0 \left(1 + \frac{\delta V_H}{\delta V_{\text{ext}}} + \frac{\delta V_{\text{xc}}}{\delta V_{\text{ext}}} \right)$$

$$\frac{\delta V_H}{\delta V_{\text{ext}}} = \frac{\delta V_H}{\delta n} \frac{\delta n}{\delta V_{\text{ext}}} = v\chi$$

$$\frac{\delta V_{\text{xc}}}{\delta V_{\text{ext}}} = \frac{\delta V_{\text{xc}}}{\delta n} \frac{\delta n}{\delta V_{\text{ext}}} = f_{\text{xc}}\chi$$

with f_{xc} = exchange-correlation kernel

Linear Response Approach

Polarizability

$$\chi \delta V_{\text{ext}} = \chi^0 (\delta V_{\text{ext}} + \delta V_H + \delta V_{\text{xc}})$$

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Linear Response Approach

Polarizability

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Linear Response Approach

Polarizability

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$$\chi = \left[1 - \chi^0 (v + f_{\text{xc}}) \right]^{-1} \chi^0$$

with f_{xc} = exchange-correlation kernel

Linear Response Approach

Polarizability

$$\chi \delta V_{\text{ext}} = \chi^0 (\delta V_{\text{ext}} + \delta V_H + \delta V_{\text{xc}})$$

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with \mathbf{f}_{xc} = **exchange-correlation kernel**

Linear Response Approach

Polarizability χ in TDDFT

① DFT ground-state calc. $\rightarrow \phi_i, \epsilon_i$ [V_{xc}]

② $\phi_i, \epsilon_i \rightarrow \chi^0 = \sum_{ij} \frac{\phi_i(\mathbf{r})\phi_j^*(\mathbf{r})\phi_i^*(\mathbf{r}')\phi_j(\mathbf{r}')}{\omega - (\epsilon_i - \epsilon_j)}$

③ $\left. \begin{aligned} \frac{\delta V_H}{\delta n} &= v \\ \frac{\delta V_{xc}}{\delta n} &= f_{xc} \end{aligned} \right\} \text{variation of the potentials}$

④ $\chi = \chi^0 + \chi^0 (v + f_{xc}) \chi$

A comment

• $f_{xc} = \left\{ \begin{aligned} &\frac{\delta V_{xc}}{\delta n} \\ &\text{"any" other function} \end{aligned} \right.$

Linear Response Approach

Polarizability χ in TDDFT

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A comment

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Finite systems

Photo-absorption cross spectrum in Linear Response

$$\sigma(\omega) = \frac{4\pi\omega}{c} \text{Im}\alpha(\omega)$$

$$\alpha(\omega) = - \int d\mathbf{r} d\mathbf{r}' V_{\text{ext}}(\mathbf{r}, \omega) \delta n(\mathbf{r}', \omega)$$

$$\sigma_{zz}(\omega) = -\frac{4\pi\omega}{c} \text{Im} \int d\mathbf{r} d\mathbf{r}' z \chi(\mathbf{r}, \mathbf{r}', \omega) z'$$

Finite systems

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Solids

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Reciprocal space

$$\chi^0(\mathbf{r}, \mathbf{r}', \omega) \longrightarrow \chi_{\mathbf{G}\mathbf{G}'}^0(\mathbf{q}, \omega)$$

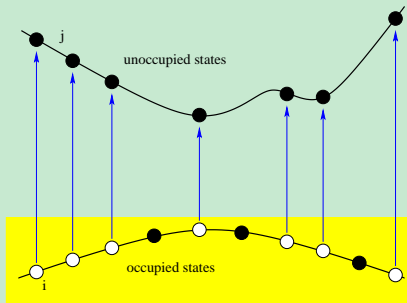
G =reciprocal lattice vector

q =momentum transfer of the perturbation

Solids

Reciprocal space

$$\chi_{\mathbf{G}\mathbf{G}'}^0(\mathbf{q}, \omega) = \sum_{\mathbf{v}\mathbf{c}\mathbf{k}} \frac{\langle \phi_{\mathbf{v}\mathbf{k}} | e^{i(\mathbf{q}+\mathbf{G})\mathbf{r}} | \phi_{\mathbf{c}\mathbf{k}+\mathbf{q}}^* \rangle \langle \phi_{\mathbf{c}\mathbf{k}+\mathbf{q}} | e^{-i(\mathbf{q}+\mathbf{G}')\mathbf{r}'} | \phi_{\mathbf{v}\mathbf{k}} \rangle}{\omega - (\epsilon_{\mathbf{c}\mathbf{k}+\mathbf{q}} - \epsilon_{\mathbf{v}\mathbf{k}}) + i\eta}$$



Solids

Reciprocal space

$$\chi_{\mathbf{G}\mathbf{G}'}^0(\mathbf{q}, \omega) = \sum_{\mathbf{v}\mathbf{k}} \frac{\langle \phi_{\mathbf{v}\mathbf{k}} | e^{i(\mathbf{q}+\mathbf{G})\mathbf{r}} | \phi_{\mathbf{c}\mathbf{k}+\mathbf{q}}^* \rangle \langle \phi_{\mathbf{c}\mathbf{k}+\mathbf{q}} | e^{-i(\mathbf{q}+\mathbf{G}')\mathbf{r}'} | \phi_{\mathbf{v}\mathbf{k}}^* \rangle}{\omega - (\epsilon_{\mathbf{c}\mathbf{k}+\mathbf{q}} - \epsilon_{\mathbf{v}\mathbf{k}}) + i\eta}$$

$$\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q}, \omega) = \chi^0 + \chi^0(\mathbf{v} + \mathbf{f}_{\text{xc}})\chi$$

$$\epsilon_{\mathbf{G}\mathbf{G}'}^{-1}(\mathbf{q}, \omega) = \delta_{\mathbf{G}\mathbf{G}'} + \mathbf{v}_{\mathbf{G}}(\mathbf{q})\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q}, \omega)$$

Solids

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$$\text{ELS}(\mathbf{q}, \omega) = -\text{Im} \{ \epsilon_{00}^{-1}(\mathbf{q}, \omega) \} ; \text{Abs}(\omega) = \lim_{\mathbf{q} \rightarrow 0} \text{Im} \left\{ \frac{1}{\epsilon_{00}^{-1}(\mathbf{q}, \omega)} \right\}$$



S.L.Adler, Phys.Rev **126**, 413 (1962); N.Wiser Phys.Rev **129**, 62 (1963)

Solids

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Reciprocal space

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Solids

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Absorption and Energy Loss Spectra $\mathbf{q} \rightarrow 0$

$$\text{ELS}(\omega) = -\text{Im} \left\{ \varepsilon_{00}^{-1}(\omega) \right\} \quad ; \quad \text{Abs}(\omega) = \text{Im} \left\{ \frac{1}{\varepsilon_{00}^{-1}(\omega)} \right\}$$

Solids

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$$\varepsilon_{00}^{-1}(\omega) = 1 + v_0 \chi_{00}(\omega)$$

Solids

Absorption and Energy Loss Spectra $\mathbf{q} \rightarrow 0$

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Solids

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Solids

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Solids

Absorption and Energy Loss Spectra $q \rightarrow 0$

$$\text{ELS}(\omega) = -\text{Im} \{ \varepsilon_{00}^{-1}(\omega) \} \quad ; \quad \text{Abs}(\omega) = \text{Im} \left\{ \frac{1}{\varepsilon_{00}^{-1}(\omega)} \right\}$$

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Exercise

$$\text{Im} \left\{ \frac{1}{\varepsilon_{00}^{-1}} \right\} = -v_0 \text{Im} \{ \bar{\chi}_{00} \}$$

Solids

Abs and ELS ($\mathbf{q} \rightarrow 0$) differs **only** by v_0

$$\text{ELS}(\omega) = -\text{Im} \{ \varepsilon_{00}^{-1}(\omega) \} \quad ; \quad \text{Abs}(\omega) = \text{Im} \left\{ \frac{1}{\varepsilon_{00}^{-1}(\omega)} \right\}$$

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$$\bar{v}_{\mathbf{G}} = \begin{cases} v_{\mathbf{G}} & \forall \mathbf{G} \neq 0 \\ 0 & \mathbf{G} = 0 \end{cases} \quad \text{microscopic components}$$

Solids

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Microscopic components \bar{v}

\bar{v} = local field effects

$$\bar{\chi}^{\text{NLF}} = \chi^0 + \chi^0 (\cancel{\chi} + f_{xc}) \bar{\chi}^{\text{NLF}}$$

Solids

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Solids

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Solids

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$$\text{Abs}^{\text{NLF}} = \text{Im} \left\{ \varepsilon_{00} \right\}$$

$$\text{Abs} = \text{Im} \left\{ \frac{1}{\varepsilon_{00}^{-1}} \right\}$$

Solids

Microscopic components \bar{v}

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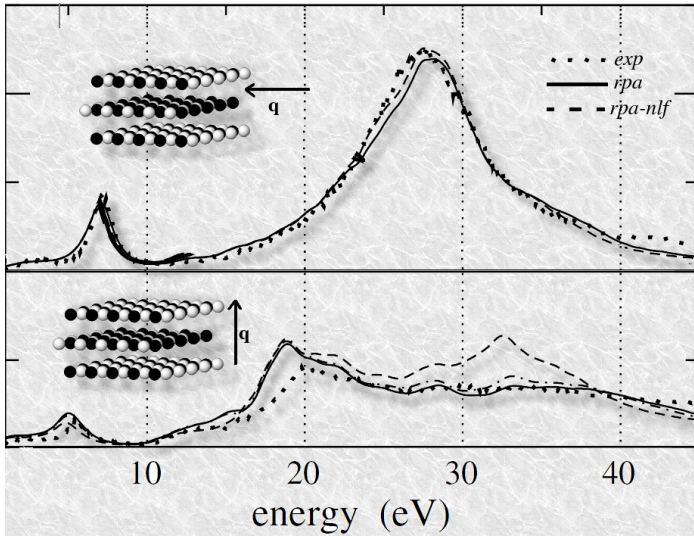
$$\text{Abs}^{\text{NLF}} = -v_0 \text{Im} \left\{ \bar{\chi}^{\text{NLF}} \right\}$$

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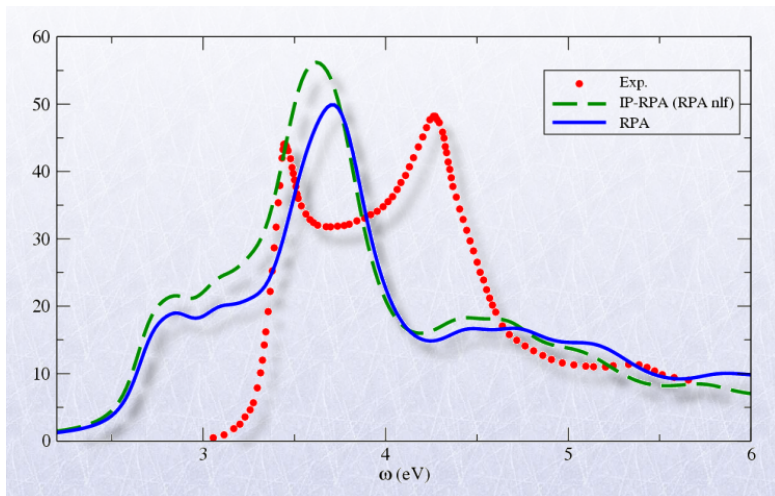
Exercise

$$\text{Abs}^{\text{NLF}} = -v_0 \text{Im} \left\{ \bar{\chi}^{\text{NLF}} \right\} = \text{Im} \{ \epsilon_{00} \}$$

EELS of Graphite



Absorption of Silicon



Outline

- 1 Introduction: why TD-DFT ?
- 2 (Just) A bit of Formalism
 - TDDFT: the Foundation
 - Linear Response Formalism
- 3 TDDFT in practice:
 - The ALDA: Achievements and Shortcomings
- 4 Resources

TDDFT in practice

Practical schema and approximations

- Ground-state calculation $\rightarrow \phi_i, \epsilon_i$ [V_{xc} LDA]
- $\chi^0(\mathbf{q}, \omega)$
- $\chi = \chi^0 + \chi^0 (v + f_{xc}) \chi$

$$f_{xc} = 0 \quad \text{RPA}$$

$$f_{xc}^{\text{ALDA}}(\mathbf{r}, \mathbf{r}') = \frac{\delta V_{xc}(\mathbf{r})}{\delta n(\mathbf{r}')} \delta(\mathbf{r} - \mathbf{r}') \quad \text{ALDA}$$

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- 4 Resources

ALDA: Achievements and Shortcomings

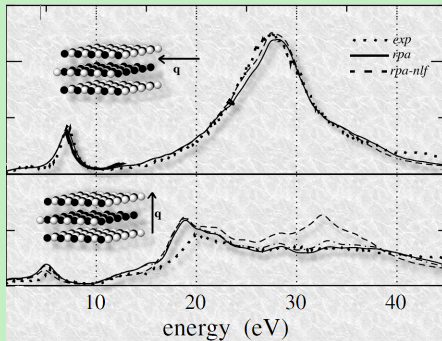
Electron Energy Loss Spectrum of Graphite

RPA vs EXP

$$\chi^{\text{NLF}} = \chi^0 + \chi^0 v_0 \chi^{\text{NLF}}$$

$$\chi = \chi^0 + \chi^0 v \chi$$

$$\text{ELS} = -v_0 \text{Im} \{ \chi_{00} \}$$



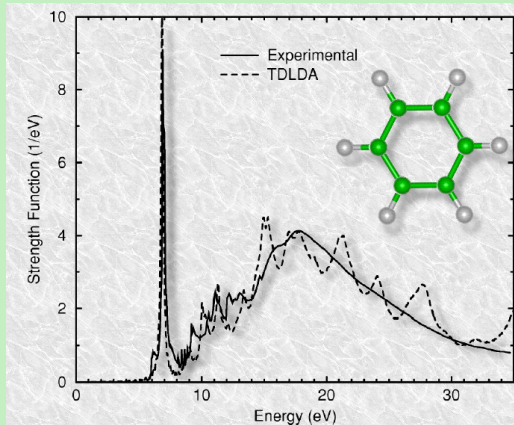
A. Marinopoulos *et al.* Phys.Rev.Lett **89**, 76402 (2002)

ALDA: Achievements and Shortcomings

Photo-absorption cross section of Benzene

ALDA vs EXP

$$\text{Abs} = -\frac{4\pi\omega}{c} \text{Im} \int d\mathbf{r} z n(\mathbf{r}, \omega)$$



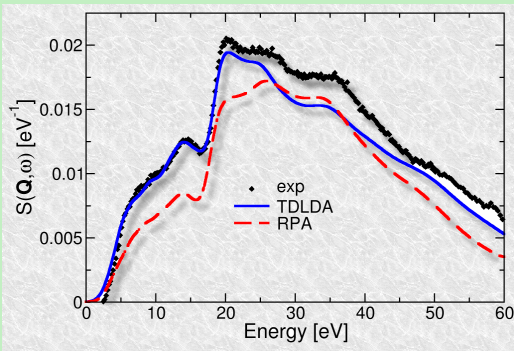
K.Yabana and G.F.Bertsch Int.J.Mod.Phys.**75**, 55 (1999)

ALDA: Achievements and Shortcomings

Inelastic X-ray scattering of Silicon

ALDA vs RPA vs EXP

$$S(\mathbf{q}, \omega) = -\frac{\hbar^2 q^2}{4\pi^2 e^2 n} \text{Im}\varepsilon_{00}^{-1}$$



Weissker *et al.*, PRL **97**, 237602 (2006)

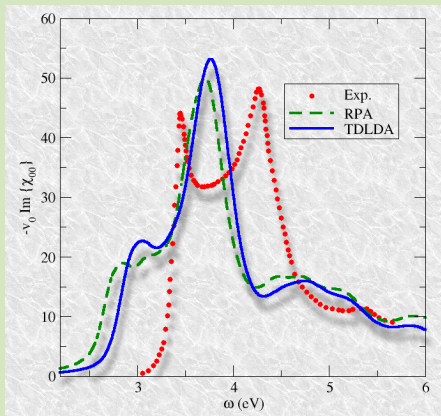
ALDA: Achievements and Shortcomings

Absorption Spectrum of Silicon

ALDA vs RPA vs EXP

$$\bar{\chi} = \chi^0 + \chi^0 (\bar{v} + f_{xc}^{ALDA}) \bar{\chi}$$

$$\text{Abs} = -v_0 \text{Im} \{ \bar{\chi}_{00} \}$$



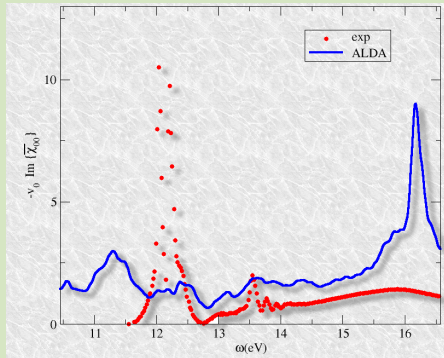
ALDA: Achievements and Shortcomings

Absorption Spectrum of Argon

ALDA vs EXP

$$\bar{\chi} = \chi^0 + \chi^0 (\bar{v} + f_{xc}^{ALDA}) \bar{\chi}$$

$$\text{Abs} = -v_0 \text{Im} \{ \bar{\chi}_{00} \}$$



ALDA: Achievements and Shortcomings

Good results

- Photo-absorption of small molecules
- ELS of solids

Bad results

- Absorption of solids

ALDA: Achievements and Shortcomings

Good results

- Photo-absorption of small molecules
- ELS of solids

Bad results

- Absorption of solids

Why?

ALDA: Achievements and Shortcomings

Good results

- Photo-absorption of small molecules
- ELS of solids

Bad results

- Absorption of solids

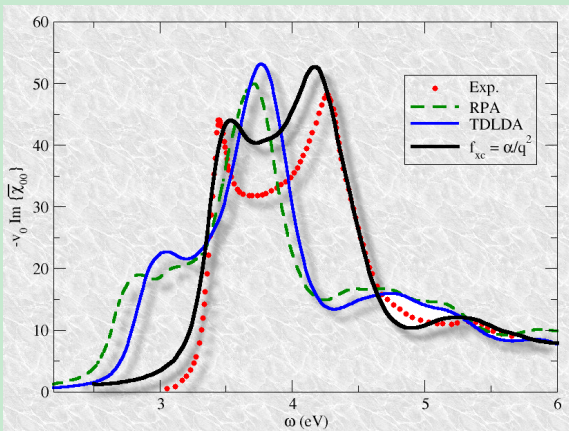
Why?

f_{xc}^{ALDA} is short-range

$$f_{xc}(\mathbf{q} \rightarrow 0) \sim \frac{1}{q^2}$$

ALDA: Achievements and Shortcomings

Absorption of Silicon $f_{xc} = \frac{\alpha}{q^2}$



L.Reining *et al.* Phys.Rev.Lett. **88**, 66404 (2002)



find better
fxc for solids

Outline

- 1 Introduction: why TD-DFT ?
- 2 (Just) A bit of Formalism
 - TDDFT: the Foundation
 - Linear Response Formalism
- 3 TDDFT in practice:
 - The ALDA: Achievements and Shortcomings
- 4 Resources

Resources

Codes (more or less) available for TDDFT

- Octopus (Marques,Castro,Rubio) -(real space, real time) - mainly finite systems - GPL
<http://www.tddft.org/programs/octopus/>
- DP (Olevano,Reining,Sottile) - (reciprocal space, frequency domain) - solides and finite systems - open source for academics
<http://dp-code.org>
- Self (Marini) - (reciprocal space, frequency domain)
- Fleszar code
- Rehr (core excitations)
- TDDFT (Bertsch)
- VASP, SIESTA, ADF, TURBOMOLE
- TD-DFPT (Baroni)

The DP code

dp - Dielectric Properties

- Reciprocal space
- Frequency domain
- Planewave basis
- Optical absorption
- Loss Spectra (EELS,IXS)

Different approximations (RPA, ALDA, NLDA, MT, etc.)

Authors: Valerio Olevano, Lucia Reining, Francesco Sottile

Valérie Vénard, Eleonora Luppi non-linear

Lucia Caramella Spin

Contributors:

Silvana Botti kernel, Wannier functions

Margherita Marsili Mapping-Theory kernel

Christine Giorgetti metals

The DP code

The algorithm

$$\varepsilon_{\mathbf{G}\mathbf{G}'}^{-1}(\mathbf{q}, \omega) = \delta_{\mathbf{G}\mathbf{G}'} + v_{\mathbf{G}}(\mathbf{q})\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q}, \omega)$$

The DP code

The algorithm

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$$\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q}, \omega) = \chi^0 + \chi^0 (v + f_{xc}) \chi$$

The DP code

The algorithm

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TDDFT vs BSE scaling

Scaling of TDDFT (DP)

$$\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q}, \omega) = \chi_{\mathbf{G}\mathbf{G}'}^0(\mathbf{q}, \omega) + \chi_{\mathbf{G}\mathbf{G}''}^0(\mathbf{q}, \omega) [v_{\mathbf{G}''}(\mathbf{q}) + f_{\mathbf{G}''\mathbf{G}'''}^{\text{xc}}(\mathbf{q}, \omega)] \chi_{\mathbf{G}'''\mathbf{G}'}(\mathbf{q}, \omega)$$

TDDFT vs BSE scaling

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N_t

TDDFT vs BSE scaling

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$N_t N_G$

TDDFT vs BSE scaling

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$$N_t N_G N_G N_\omega$$

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Scaling of TDDFT (DP)

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$$N_t N_G N_G N_\omega \quad ; \quad N_t N_r \ln N_r N_\omega$$

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$$N_t N_G N_G N_\omega \quad ; \quad N_t N_r \ln N_r N_\omega$$

$$\text{Scaling} \lesssim N_{at}^4$$

TDDFT vs BSE scaling

Scaling of BSE (EXC)

$$H_{vck}^{v'c'k'}$$

Scaling $N_t^2 N_r \lesssim N_{at}^5$

TDDFT vs BSE scaling

Scaling of BSE (EXC)

$$H_{vck}^{v'c'k'} A_{\lambda}^{v'c'k'} = A_{\lambda} A_{\lambda}^{vck}$$

$$\text{Scaling} \quad N_t^2 N_r \lesssim N_{at}^5$$

$$\text{Scaling} \quad N_t^3 \lesssim N_{at}^6$$

TDDFT

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oooooooo

Some References

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